

Gas-phase CO₂ emission toward Cepheus A East: the result of shock activity?¹

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ABSTRACT

We report the first detection of gas-phase CO₂ emission in the star-forming region Cepheus A East, obtained by spectral line mapping of the ν_2 bending mode at 14.98 μm with the Infrared Spectrograph (IRS) instrument onboard the *Spitzer Space Telescope*. The gaseous CO₂ emission covers a region about $35'' \times 25''$ in extent, and results from radiative pumping by 15 μm continuum photons emanating predominantly from the HW2 protostellar region. The gaseous CO₂ exhibits a temperature distribution ranging from 50 K to 200 K. A correlation between the gas-phase CO₂ distribution and that of H₂ S(2), a tracer of shock activity, indicates that the CO₂ molecules originate in a cool post-shock gas component associated with the outflow powered by HW2. The presence of CO₂ ice absorption features at 15.20 μm toward this region and the lack of correlation between the IR continuum emission and the CO₂ gas emission distribution further suggest that the gaseous CO₂ molecules are mainly sputtered off grain mantles – by the passage of slow non-dissociative shocks with velocities of 15–30 km s^{−1} – rather than sublimated through grain heating.

Subject headings: ISM: Molecules — ISM: Clouds, molecular processes — Star-forming region: individual (Cepheus A East)

¹Based on observations with the *Spitzer Space Telescope*.

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1. Introduction

Cepheus A is a well-known site of star formation that has been studied extensively. Previous results have shown that it contains a series of heavily embedded far-infrared and radio-continuum sources, one of which dominates the luminosity of the entire region (HW2 with $2.5 \times 10^4 L_{\odot}$; Hughes & Wouterloot 1984; Evans et al. 1981). Ground- (e.g. Goetz et al. 1998) and space-based observations (e.g. Hartigan et al. 1996; Wright et al. 1996; van den Ancker et al. 2000) revealed that the multipolar outflow existing in this region consists of complex structures of shock-excited atomic and molecular gas. H_2 emissions and line emissions from ionized species were detected along high-velocity (HV) and extra-high-velocity (EHV) outflows, indicating the presence of both dissociative (J-type) and non-dissociative (C-type) shocks, likely resulting from successive episodes of activity (Narayanan & Walker 1996). The protostellar object HW2 was determined to be the dominant powering source of the EHV outflow, while the source of the HV jet is still under debate (Goetz et al. 1998; Hiriart et al. 2004). The cavities carved into the surrounding molecular gas by the quadrupolar outflow are clearly seen in the NH_3 maps obtained by Torrelles et al. (1993).

In this *Letter*, we focus on new observations of Cepheus A East obtained with the Infrared Spectrograph (IRS) onboard the *Spitzer Space Telescope* (*SST*). The benefit of the *SST* data lies in the much greater sensitivity and the significantly better spatial resolution offered by the IRS compared to the *Infrared Space Observatory* (*ISO*). Such advances allow for the detection and spectral mapping of variations in the molecular and atomic line intensities on scales of a few arcsec instead of the previous few tens of arcsec or few arcmin scales, thus enabling our understanding of shock physics and chemistry at much finer scales. Detections of H_2 S(1) through S(7), C_2H_2 , [Ne II], [Ne III], [S I], [S III] and [Fe II] emissions as well as absorption from CO_2 and H_2O ices will be discussed in a future paper. The present work will focus on the first detection of gas-phase CO_2 emission toward Cepheus A East. Section 2 describes the observations and data analysis. Sections 3 & 4 discuss our new results as well as the spatial distributions of gas-phase CO_2 and H_2 S(2) in the context of shock chemistry and outflow activity.

2. Observations and data reduction

Spectral maps of two overlapping $1' \times 1'$ -square fields were obtained toward Cepheus A East with the IRS instrument onboard the *Spitzer Space Telescope* (*SST*) as part of Guaranteed Time Observer (GTO) program 113. The Short-Low (SL, both orders), Short-High (SH) and Long-High (LH) modules allowed for wavelength coverage from 5.2 to 25 μm . Most of the data longward of 25 μm suffer severe detector saturation as a result of

very strong continuum emission and were not used. Continuous spatial coverage in the overlapping fields was obtained by stepping the slit perpendicular and parallel to its long axis in steps of one-half its width and 4/5 its length, respectively.

The data were processed at the Spitzer Science Center (SSC) using version 12 of the IRS pipeline. The Spectroscopy Modeling Analysis and Reduction Tool (SMART) software (Higdon et al. 2004) was used to extract wavelength-calibrated spectra at each spatial position in the maps. We used a set of locally-developed IDL routines specifically designed to remove bad pixels in the high resolution modules, to calibrate the fluxes for extended sources, and to generate spectral line maps from the extracted spectra (Neufeld et al. 2006).

3. Results

Figure 1 shows two examples of spectral line intensity maps we obtained using the SH data alone to allow direct comparison of the H₂ S(2) and gaseous CO₂ spatial distributions. The right panel displays the spatial distribution of the H₂ S(2) intensity (12.28 μ m), while the left panel shows the intensity distribution for gas-phase CO₂ (*Q*-branch of the ν_2 bending mode at 14.98 μ m). The H₂ and the CO₂ emissions mainly arise at the surfaces of the NH₃-free cavities carved by the outflows. These emissions thus seem to result from interactions between the EHV outflow and the quiescent molecular gas traced by NH₃ (Gómez et al. 1999). Figure 2 (left panel) displays examples of summed spectra in the wavelength range relevant for this study.

To constrain the physical conditions in the CO₂-containing gas, we generated synthetic profiles of the ν_2 CO₂ band for temperatures ranging from 50 K to 900 K and compared them with the observed spectra (Fig. 2, right panel). We find that temperatures between 50 and 200 K best fit the observed CO₂ *Q*-branch emission (14.98 μ m) over the entire CO₂-emitting region. Our observations, therefore, reveal the presence of a post-shock gas component much cooler than that measured with H₂ pure rotational line transitions (e.g. van den Ancker et al. 2000, $T \sim 730$ K from H₂ S(1)–S(5) over the *ISO* Short Wavelength Spectrometer beam sizes 14'' \times 20'' and 14'' \times 27'').

To determine the column density associated with the gas-phase CO₂ emission, one needs to identify the excitation mechanism. Prior studies (e.g., González-Alfonso & Cernicharo 1999; Boonman et al. 2003) showed that gaseous CO₂ molecules can be excited: 1) by collisions with H and H₂; 2) by radiative pumping into the 4.27 μ m band; 3) by radiative pumping due to 15 μ m continuum photons emitted by dust local either to the CO₂ gas component or to HW2.

For collisions to dominate, densities in excess of $n \sim 10^8 \text{ cm}^{-3}$ would be required. However, prior observations indicate densities from $\sim 10^3$ to $\sim 10^7 \text{ cm}^{-3}$ in the outflow region (Codella et al. 2005). Radiative pumping into the $4.27 \mu\text{m}$ band and subsequent cascade can be disregarded because the relaxation to the $15 \mu\text{m}$ band would produce detectable emission at $13.9 \mu\text{m}$ and $16.2 \mu\text{m}$ that we do not observe. Resonant scattering by $15 \mu\text{m}$ continuum photons is, therefore, the most likely mechanism to excite the observed gaseous CO_2 molecules. Continuum maps obtained from our data also indicate that the radiation field is dominated by dust emission close to HW2 – from a reflection nebula (Casement & McLean 1996; Goetz et al. 1998, Martín-Pintado et al. 2005) – rather than by dust emission local to the CO_2 gas component. We, hence, conclude that the gas-phase CO_2 molecules detected toward Cepheus A East are predominantly excited via radiative pumping by $15 \mu\text{m}$ continuum photons emanating from the protostellar region HW2. Resonant scattering in the CO and H_2O vibrational bands, through continuum photons emanating from IRc2/BN, was also found toward the shock region Orion Peak 1/2 (González-Alfonso et al. 2002).

To compute $N(\text{CO}_2)$, we adopted the total luminosity and source effective temperature from Lenzen et al. (1984), as well as a distance of 690 pc for the Cepheus A East region, and we assumed that the projected distance of each spatial position is equal to its true distance from HW2. Under that assumption, the inferred CO_2 column density is proportional to the measured intensity and to the square of the angular separation from the exciting source HW2. We also corrected the CO_2 intensity measurements for extinction caused by the CO_2 ice absorption. A comparison of the $N(\text{CO}_2)$ map (Fig. 3, left panel) with the intensity distribution of the H_2 S(2) line ($I(\text{H}_2)$ S(2), Fig. 1, right panel) indicates that – like H_2 S(2) – CO_2 peaks at the NE position, coincident with the NH_3 bridge between Cep-A2 and Cep-A3 (Torrelles et al. 1993) and at the HW5/HW6 positions. These striking similarities again strongly suggest that the CO_2 emission arises in post-shock gas, traced by the warm H_2 emission, and results from interactions between the EHV outflow and the ambient medium. The right panel of Fig. 3 displays the derived, extinction corrected, column density measurements of gas-phase CO_2 *versus* H_2 S(2) for spatial positions where the CO_2 line intensity (Q -branch) is greater than $0.75 \times 10^{-4} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ sr}^{-1}$. Our estimate of the Spearman coefficient ($\rho \sim 0.42$ and significance level $< 1\%$) indicates that the column densities of the two species tend to increase together with spatial position.

4. Discussion

The presence of CO_2 ice absorption features at most spatial positions in the gaseous CO_2 -emitting region, their progressive weakening along the EHV outflow axis, as well as

the localization of the gaseous CO_2 emission into a cool post-shock component, all pose the question of the origin of the gas-phase CO_2 molecules in Cepheus A East. Was CO_2 injected into the gas-phase by sublimation from CO_2 -rich dust mantles after substantial grain heating or by grain sputtering caused by the passage of slow non-dissociative shocks known to exist in this active region?

We searched for evidence of a correlation between the dust continuum emission, a tracer of grain heating, and the gaseous CO_2 distribution in the CO_2 -emitting region by comparing the CO_2 column density measurements with measurements of the dust continuum emission at $7.3\ \mu\text{m}$, $14.8\ \mu\text{m}$ and $18.0\ \mu\text{m}$. We found no obvious relation between $N(\text{CO}_2)$ and the dust continuum emission at these wavelengths. Furthermore, we determined that the dust temperature is never greater than the CO_2 sublimation temperature ($\sim 90\ \text{K}$) at the locations of the CO_2 column peaks. While grain heating likely occurs in the active region close to HW2, CO_2 -rich grain mantle sublimation is not the dominant source of gaseous CO_2 in Cepheus A East. Sputtering off CO_2 -rich ice mantles is, therefore, the most likely scenario in the present case. Our measurements of the CO_2 ice column density range from $\sim 10^{17}$ to $\sim \text{few} \times 10^{18}\ \text{cm}^{-2}$ over the region exhibiting CO_2 gas emission, indicating that only a few percent of the material has passed through mantle-destroying shocks.

To derive the abundance of gaseous CO_2 , an estimate of $N(\text{H}_2)$ in the CO_2 gas component is needed. Because the gas-phase CO_2 molecules arise from a post-shock component much cooler than that containing the H_2 gas we detected, we cannot directly measure the total column density of the gas at $T = 50\text{--}200\ \text{K}$ from our data. Our abundance estimate will, hence, rely on previous indirect measures. Gómez et al. (1999) derived $N(\text{H}_2) = 1.5 \times 10^{22}\ \text{cm}^{-2}$ in the outflow using HCO^+ emission measurements; a similar estimate was obtained by van den Ancker (2000) using far-infrared CO emission at the *ISO*/Long Wavelength Spectrometer angular resolution ($75''$ beam size). Adopting this H_2 column density yields an average gas-phase CO_2 abundance of $\text{few} \times 10^{-7}$.

This value should be regarded as a lower limit since the $N(\text{CO}_2)$ might be underestimated if the true distance from HW2 is greater than the projected distance we assumed. A small fraction of the gaseous CO_2 might arise from CO_2 ice sublimation since some CO_2 ice profiles, especially close to HW2, show substructures characteristic of grain heating (Gerakines et al. 1999). A fraction of the gas-phase CO_2 emission might also come from quiescent gas associated with Cep-A2 (knot at [R.A., Dec] = $[7'', 18'']$ in Fig. 1). While these additional sources of gaseous CO_2 certainly contribute to the scatter seen in the right panel of Fig. 3, none of these sources can account for the correlation we observe between $N(\text{CO}_2)$ and $N(\text{H}_2)$ S(2) over such a large spatial extent.

The low gas-phase CO_2 abundances detected by *ISO* ($\sim 10^{-7}$ instead of the expected

few $\times 10^{-6}$) toward active star-forming regions with large CO₂ ices reservoirs (e.g. van Dishoeck et al. 1996) fostered theoretical efforts to investigate the effects that shocks might have on CO₂ chemistry. Charnley & Kaufman (2000) showed that for shocks faster than a critical shock speed, which is a function of preshock gas density, CO₂ molecules sputtered off grain mantles are expected to be efficiently destroyed in the shock through reactions with H and H₂. On the other hand, a shock speed higher than 15 km s⁻¹ is necessary to allow efficient sputtering.

Comparison of our results with the Charnley & Kaufman (2000) shock-model predictions indicates that slow C-type shocks with speeds of at least ~ 15 km s⁻¹ interacted with the ambient gas over the extent of the CO₂-emitting region (see Fig 1). Adopting a pre-shock density of $n_H \sim 10^5$ cm⁻³ (Goetz et al. 1998), our abundance limit further indicates that the shock speed is no greater than 30 km s⁻¹ since, at this density, higher shock speeds would efficiently remove CO₂ from the gas phase, mainly through reactions with hydrogen atoms (see Fig. 3 of Charnley & Kaufman 2000). The CO₂ emission might arise at the flanks of the bow shock modeled by Froeblich et al. (2002). The average gas phase CO₂ abundance derived above (few $\times 10^{-7}$) is consistent with a scenario in which a few percent of the grain material has been subject to shocks fast enough to sputter CO₂ ice mantles (with CO₂(ice)/H₂ \sim few $\times 10^{-5}$).

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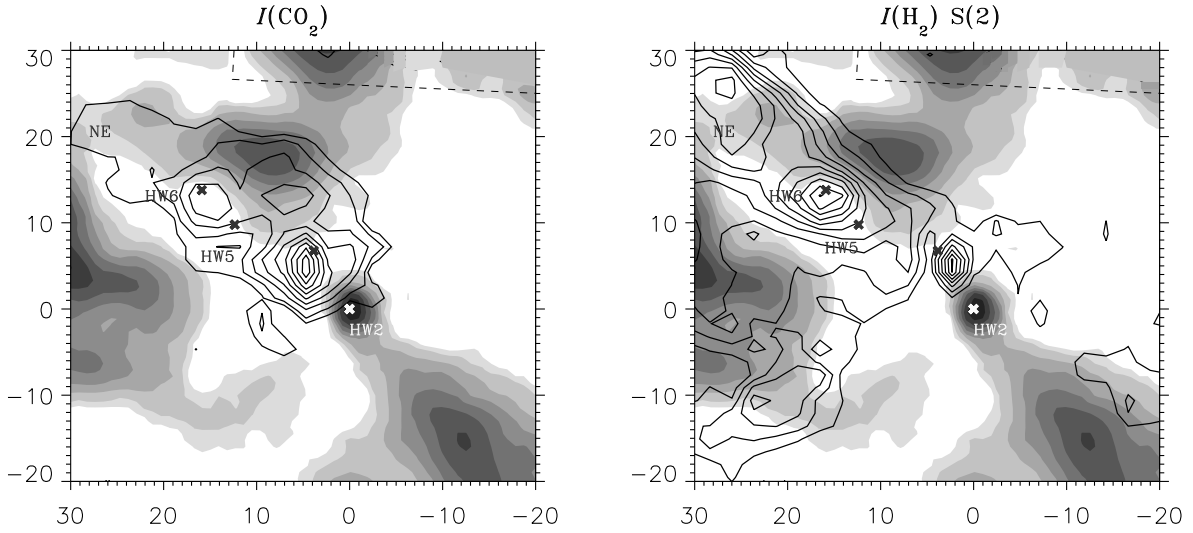


Fig. 1.— Intensity distribution of gas-phase CO_2 emission (Q -branch at $14.98 \mu\text{m}$) and H_2 S(2) pure rotational line emission (at $12.28 \mu\text{m}$) toward Cepheus A East is shown as black contours. The lowest intensity contours are 0.75 and 1.0 for CO_2 and 1.0 for H_2 S(2) with steps of $0.5 \times 10^{-4} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ sr}^{-1}$. The (x) indicate the position of the radio-continuum sources HW2, HW4, HW5, HW6 and the EHV outflow edge (NE: $\Delta\alpha \cos \delta = 27''$; $\Delta\delta = 22''$). The coordinates are offsets in R.A. ($\Delta\alpha \cos \delta$) and declination ($\Delta\delta$) in arcsec with respect to HW2 (J2000: $\alpha = 22\text{h}56\text{m}17\text{s}.9$ and $\delta = +62^\circ 01' 49''$; Hughes & Wouterloot 1984). The grayscale shows the distribution of NH_3 (1,1), a tracer of the cold quiescent molecular gas. The lowest contours are 10 and 25 with steps of $25 \text{ mJy km s}^{-1} \text{ beam}^{-1}$ (Torrelles et al. 1993). The black dashed line delineates the mapped region in Short-High.

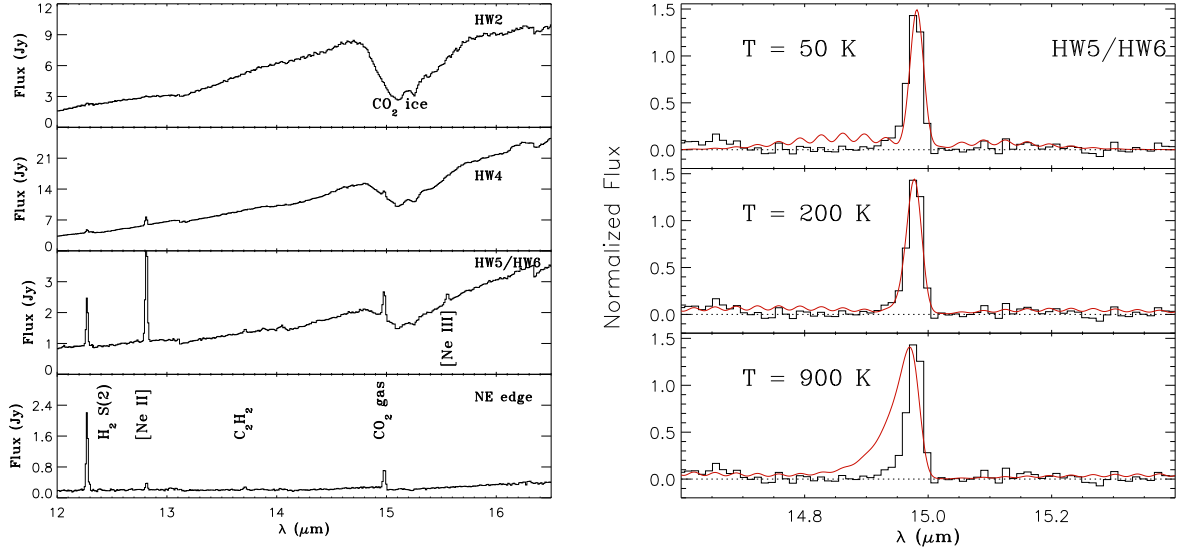


Fig. 2.— *left panel*: Summed IRS spectra of gas-phase CO_2 , H_2 , C_2H_2 , $[\text{Ne II}]$ and $[\text{Ne III}]$ emission and CO_2 ice absorption toward Cepheus A East. Individual spectra were summed over regions $\sim 6'' \times 8''$ in size, centered on the radio-continuum sources indicated in Fig. 1. Note the appearance of the gas-phase CO_2 emission feature and the progressive decrease in the CO_2 ice absorption feature when moving along the axis of the EHV outflow traced by the H_2 emission and the NH_3 cavities and away from HW2. Note also the presence of gaseous CO_2 emission alone at the NE position, the farthest from HW2 in our maps. *Right panel*: Continuum-subtracted summed spectrum from the HW5/HW6 position. Overimposed (red trace) is the calculated CO_2 band spectrum for $T = 50, 200$ and 900K . Note the progressive shift of the bandhead and the Q -branch widening with increasing temperature.

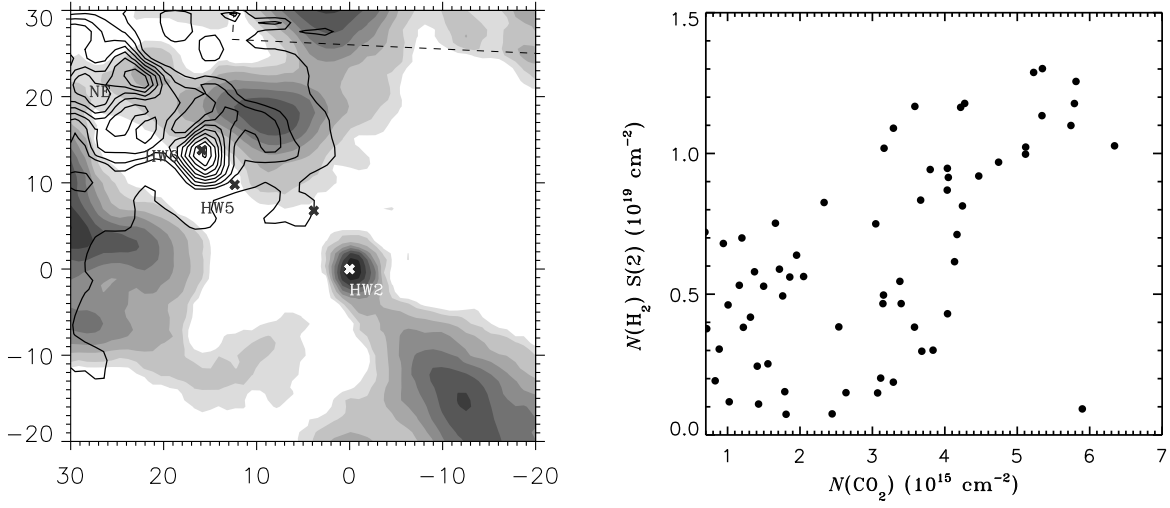


Fig. 3.— *Left panel:* Distribution of gas-phase CO₂ column density, $N(\text{CO}_2)$. The lowest contours are 1.5 and 3 with steps of $0.5 \times 10^{15} \text{ cm}^{-2}$. The symbols and designations are as in Fig. 1. *Right panel:* $N(\text{CO}_2)$ vs $N(\text{H}_2) J = 4$ (S(2) line) for the spatial positions exhibiting gaseous CO₂ (Q -branch) intensities starting at $0.75 \times 10^{-4} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ sr}^{-1}$.